



## **Revising the WASP-Park's model wake decay coefficient for different atmospheric stabilities to better predict wind farm outputs**

**Pena Diaz, Alfredo; Rathmann, Ole; Réthoré, Pierre-Elouan**

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## Introduction

Prediction of wind farm power outputs is difficult, particularly for large-size farms and when the variation in stability conditions is strong. CFD methods are applied for multiple turbines using LES and on larger wind farms using RANS turbulence models. LESs remain too expensive and still have problems to preserve turbulence over large distances. RANS methods have numerical and physical limitations and approach the multi-scale physics of turbine wake using a simplistic single scale model [6]. The majority of the CFD methods assume neutral atmospheric conditions. Although engineering wake models are based in many assumptions and limited to medium-size wind farms, they are million times faster than RANS/LES and can be used for short-term farm power predictions under a wide range of scenarios. The Park model [3] used in WAsP [4] makes use of the wake decay coefficient  $k_w$  to estimate the speed deficit downstream of a(a cluster of) wind turbine(s). [5] found that in order to match the speed deficits estimated by the infinite wind farm boundary layer (IWFB) model of Frandsen [2] to those of the Park model, the wake decay had to be modified depending on stability, roughness and turbine-turbine distance. Here we present part of the findings from [5] and discuss their main implications

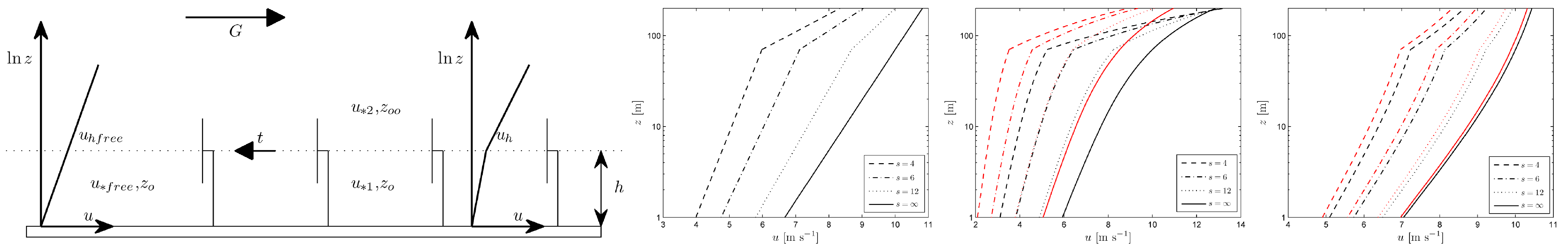


Figure 1: (left) The infinite wind farm boundary layer model of Frandsen. Offshore wind profiles for different stabilities: neutral- $L = \infty$  m (middle), stable- $L = 200$  m (right), unstable- $L = -200$  m (right most). Versions 1 (black) and 2 (red) with  $C_t = 0.88$ ,  $h = 70$  m and  $u_{hfree} = 10$  m s $^{-1}$

## Background

The IWFB model of Frandsen assumes a two-layer atmosphere within an infinite wind farm (above and below hub height  $h$ —Fig. 1). By linking the shear of both layers at  $h$ , he found the wind speed reduction  $R_u$ , i.e. the ratio of the wind speed within the farm  $u_h$  to the free wind speed,

$$R_u = \frac{1 + K_2 \sqrt{K_1^{-2}}}{1 + K_2 \sqrt{K_1^{-2} + c_t}} \quad (1)$$

where  $K_2 = (1/\kappa) \ln[G/(h f_p)]$ ,  $K_1 = (1/\kappa) \ln(h/z_o)$ ,  $f_p = f_c \exp(A_*)$ ,  $c_t = (\pi/8) [C_t/(s_r s_f)]$ ,  $\kappa$  being the von Kármán constant,  $G$  the geostrophic wind speed,  $z_o$  the roughness length,  $f_c$  the Coriolis parameter,  $A_*$  a modified  $A$ -parameter from the resistant law constants,  $C_t$  the thrust coefficient, and  $s_r$  and  $s_f$  the along and cross dimensionless turbine-turbine distances, respectively.

Following the approach by Emeis [1] one can extend Eq. 1 by multiplying the dimensionless wind shear  $\phi_m$  from MOST to  $K_2$  and subtracting the stability correction  $\psi_m$  to the log term in  $K_1$ . Both  $\phi_m$  and  $\psi_m$  are functions of height and of the Obukhov length  $L$ . We called this version 1.

Another approach (hereafter version 2) is to extend the shear stresses  $u_{*1}$  and  $u_{*2}$  using the  $\psi_m$  functions and the diabatic wind profile. Following the derivation of Eq. (1),  $K_1$  has to be corrected as in version 1 and  $K_2$  becomes  $u_*$  dependent, since  $A_*$  is  $u_*$  dependent through  $\mu_o$  ( $\mu_o = \kappa u_*/(f_c L)$ ). Therefore,  $K_2$  is different for the undisturbed atmosphere (in the numerator in Eq. 1 we use  $K_2(u_{*free})$ ) than for the atmosphere within the farm (in the denominator in Eq. 1 we use  $K_2(u_{*2})$ ).  $u_{*2}$  is found iteratively using the expression for the effective roughness  $z_{oo}$ ,

$$z_{oo} = h \exp \left[ -\kappa / \sqrt{c_t + K_1^{-2}} - \psi_m(h/L) \right] \quad (2)$$

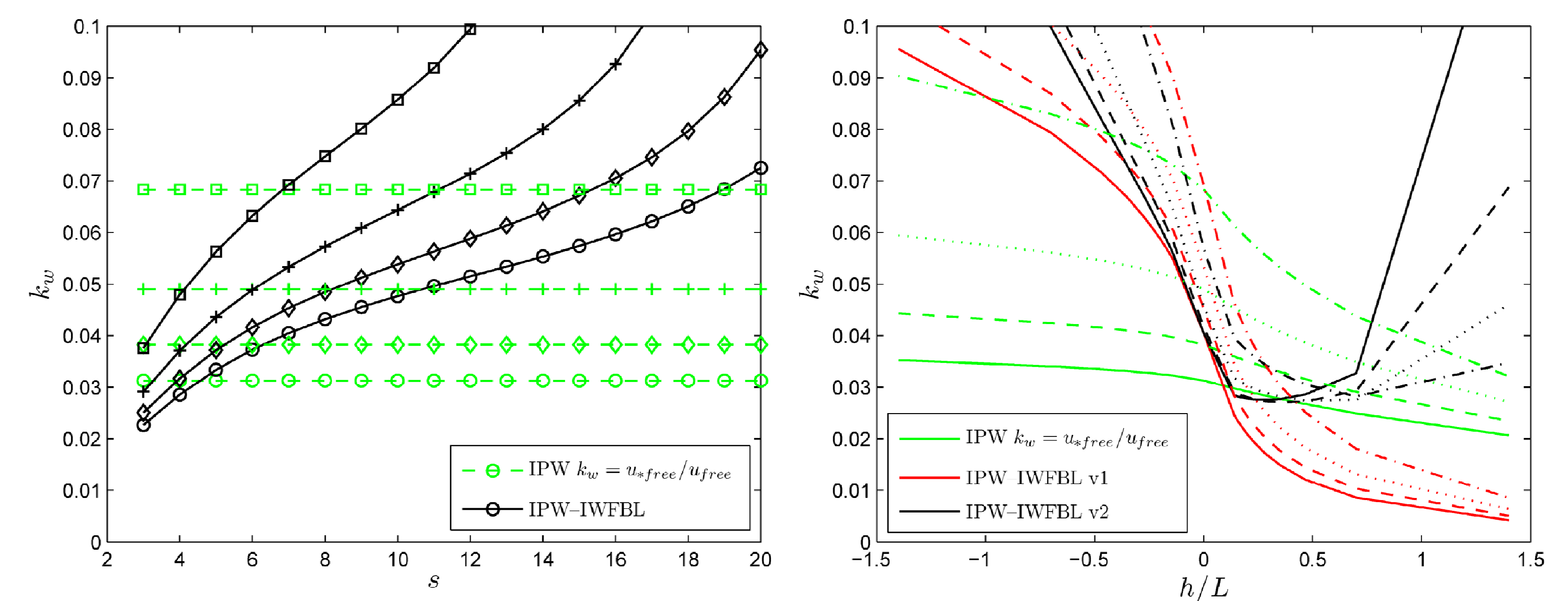


Figure 3: Wake decay coefficient for neutral (left) and diabatic conditions (right) with  $s = 7$ ,  $C_t = 0.88$ ,  $h = 70$  m and  $u_{hfree} = 10$  m s $^{-1}$

## Results &amp; Discussion

Fig. 1 shows wind profiles within an offshore farm for different turbine distances and stabilities, where the kink in the profile appears at  $h$ . Both versions show the same wind profiles under neutral conditions and for the unstable and stable cases, version 1' wind speeds are lower than those of version 2.

Fig. 2-left shows  $R_u$  for a range of stabilities and roughnesses. For version 1,  $R_u$  always increases with  $z_o$  with a higher stability variation. Version 2 follows the behavior of 1 with a peak at a positive  $h/L$ . This is because  $u_{hfree}$  is used as the input parameter and for  $h = 70$  m and stable conditions, the stability corrections from MOST to the wind speed are too high. The peak is also observed in the normalized geostrophic wind speed graph in Fig. 2-right.

The wake decay coefficient has been estimated for the infinite Park model so that it matches  $R_u$  of the two versions of the extended IWFB model. In Fig. 3 this adjusted  $k_w$  is compared to the expression  $k_w = u_{*free}/u_{hfree}$  for neutral (left frame) and diabatic conditions (right frame). For a wide number of turbine distances, roughnesses and stabilities, the adjusted  $k_w$  is much lower than the recommended WAsP value (0.050 and 0.075 for offshore and onshore sites, respectively). The value of the adjusted  $k_w$  increases with increasing roughness, turbine separation and instability.

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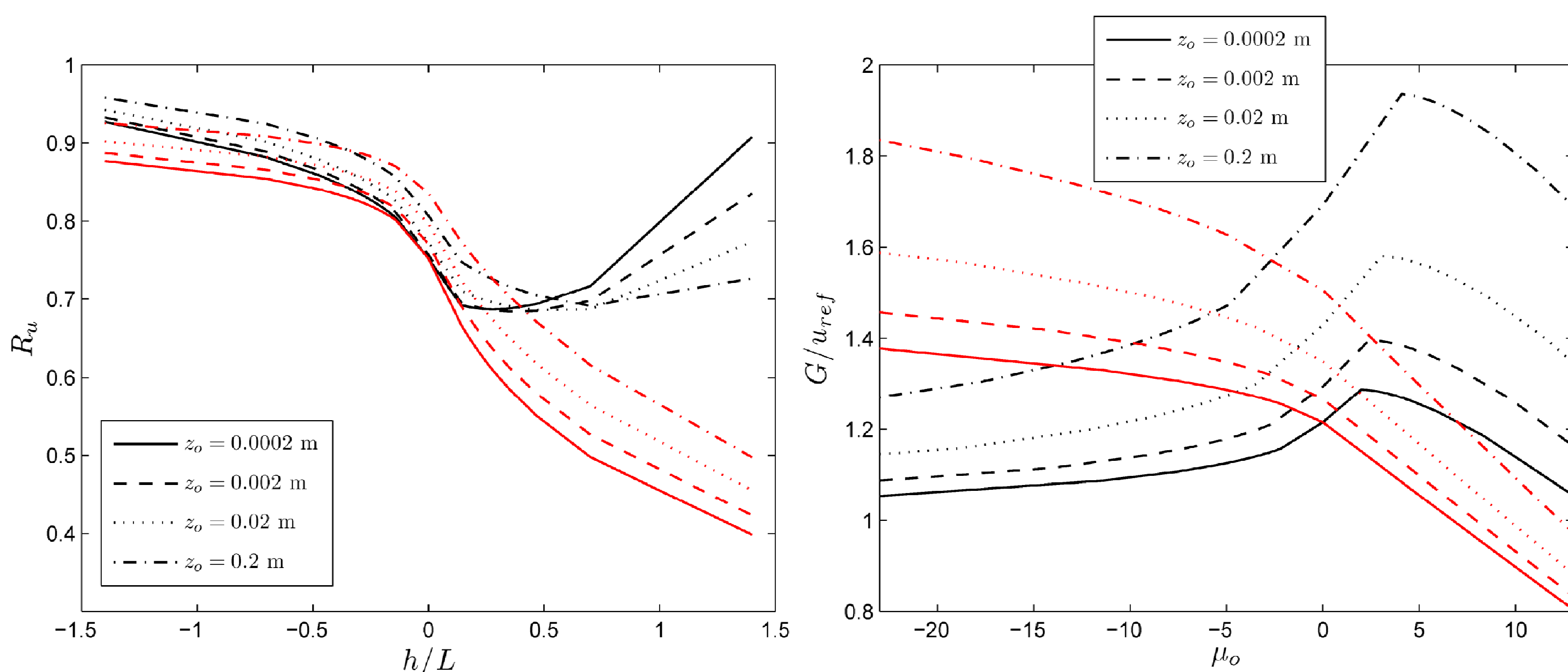


Figure 2: Wind speed reduction (left) and normalized geostrophic wind speed (right) for versions 1 (red) and 2 (black) with  $s = 7$ ,  $C_t = 0.88$ ,  $h = 70$  m and  $u_{hfree} = 10$  m s $^{-1}$